

Figure 3.3: Depiction of a train consisting of cars treated as masses $M$ and linkages treated as springs of stiffness $K$ or compliance $C=1 / K$. Below it is the electrical equivalent circuit for comparison. The masses are modeled as inductors and the springs as capacitors to ground. The velocity is analogous to a current and the force $f_{n}(t)$ to the voltage $\phi_{n}(t)$. The length of each cell is $\Delta$ [ m ]. The train may be accurately modeled as a transmission line (TL), since the equivalent electrical circuit is a lumped model of a TL. This method, called a Cauer synthesis, is based on the ABCD transmission line method of Sec. 3.8 (p. 105).

### 3.3.3 Transmission-line analysis

Problem \# 2: Train-mission-line We wish to model the dynamics of a freight train that has $N$ such cars and study the velocity transfer function under various load conditions.

As shown in Fig. 4.11, the train model consists of masses connected by springs.
Use the ABCD method (see the discussion in Appendix B.3, p. 228) to find the matrix representation of the system of Fig. 4.11. Define the force on the $n$th train $\operatorname{car} f_{n}(t) \leftrightarrow F_{n}(\omega)$ and the velocity $v_{n}(t) \leftrightarrow V_{n}(\omega)$.

Break the model into cells consisting of three elements: a series inductor representing half the mass ( $M / 2$ ), a shunt capacitor representing the spring $(C=1 / K)$, and another series inductor representing half the mass ( $L=M / 2$ ), transforming the model into a cascade of symmetric $(\mathcal{A}=\mathcal{D})$ identical cell matrices $\mathcal{T}(s)$.

- 2.1: Find the elements of the $A B C D$ matrix $\mathcal{T}$ for the single cell that relate the input node 1 to output node 2

$$
\left[\begin{array}{c}
F  \tag{DE-3.3}\\
V
\end{array}\right]_{1}=\mathcal{T}\left[\begin{array}{c}
F(\omega) \\
-V(\omega)
\end{array}\right]_{2}
$$

Sol:

$$
\begin{align*}
\mathcal{T} & =\left[\begin{array}{cc}
1 & s M / 2 \\
0 & 1
\end{array}\right]\left[\begin{array}{cc}
1 & 0 \\
s C & 1
\end{array}\right]\left[\begin{array}{cc}
1 & s M / 2 \\
0 & 1
\end{array}\right] \\
& =\left[\begin{array}{cc}
1+s^{2} M C / 2 & (s M)\left(1+s^{2} M C / 4\right) \\
s C & 1+s^{2} M C / 2
\end{array}\right] \tag{DE-3.4a}
\end{align*}
$$

$\square$

- 2.2: Express each element of $\mathcal{T}(s)$ in terms of the complex Nyquist ratio $s / s_{c}<1$ ( $s=2 \pi j f, s_{c}=2 \pi j f_{c}$ ). The Nyquist wavelength sampling condition is $\lambda_{c}>2 \Delta$. It says the critical wavelength $\lambda_{c}>2 \Delta$. Namely it is defined in terms the minimum number of cells $2 \Delta$, per minimum wavelength $\lambda_{c}$.
The Nyquist wavelength sampling theorem says that there are at least two cars per wavelength.
Proof: From the figure, the distance between cars $\Delta=c_{o} T_{o}[\mathrm{~m}]$, where

$$
c_{o}=\frac{1}{\sqrt{M C}} \quad[\mathrm{~m} / \mathrm{s}]
$$

The cutoff frequency obeys $f_{c} \lambda_{c}=c_{o}$. The Nyquist critical wavelength is $\lambda_{c}=c_{o} / f_{c}>2 \Delta$. Therefore the Nyquist sampling condition is

$$
\begin{equation*}
f<f_{c} \equiv \frac{c_{o}}{\lambda_{c}}=\frac{c_{o}}{2 \Delta}=\frac{1}{2 \Delta \sqrt{M C}} \quad[\mathrm{rad} / \mathrm{sec}] \tag{DE-3.5}
\end{equation*}
$$

Finally, $s_{c}=\jmath 2 \pi f_{c}$.

Sol: The solution is a repeat what is summarized above: the system in Fig. 4.11 represents a transmission line having a wave speed of $c_{o}=1 / \sqrt{M C}$ and characteristic impedance $r_{o}=\sqrt{M / C}$. Each cell, composed of 2 masses $M$ connected by one spring $K$, has length $\Delta$.

We wish to define the Nyquist frequency $f_{c}$ such that the wavelength $\lambda>2 \Delta$, where $\Delta$ is the cell length. Using the formula for the wavelength in terms of the wave velocity and frequency we find

$$
\lambda=c_{o} / f_{c}=2 \Delta
$$

thus we conclude that

$$
\begin{equation*}
f<f_{c}=\frac{c_{o}}{2 \Delta}=\frac{1}{2 \Delta \sqrt{M C}} \tag{DE-3.6}
\end{equation*}
$$

If we wish to have the system be accurate for a given frequency we may make the cell length $\Delta$ smaller, while keeping the velocity constant ( $M C$ is held constant). Thus the characteristic resistance [ohms/unit length] $r_{o}$ must change as $f_{c} \rightarrow \infty$ and $\Delta \rightarrow 0$. We can either let $M \rightarrow \infty$ and $C \rightarrow 0$ (their product remains constant), or the other way around. In one case $r_{o} \rightarrow \infty$ and in the other case it goes to 0 .

- 2.3: Use the property of the Nyquist sampling frequency $\omega<\omega_{c}(E q . D E-3.4)$ to remove higher order powers of frequency

$$
\begin{equation*}
1+\left(\frac{s}{s_{c}}\right)^{2^{2}} \approx 1 \tag{DE-3.7}
\end{equation*}
$$

to determine a band-limited approximation of $\mathcal{T}(s)$.
Sol:

$$
\begin{aligned}
\mathcal{T} & =\left[\begin{array}{cc}
1+2\left(s / s_{c}\right)^{2} & s M\left(1+\left(s / s_{c}\right)^{2}\right) \\
s C & 1+2\left(s / s_{c}\right)^{2}
\end{array}\right] \\
& \approx\left[\begin{array}{cc}
1 & s M \\
s C & 1
\end{array}\right]
\end{aligned}
$$

The approximation is highly accurate below the Nyquist cutoff frequency $s<s_{c}$. Given any desired frequency $f$, we can always make the cell size $\Delta$ smaller by decreasing $M$ and $C$, while keeping $f<f_{c}$ and the cell velocity constant ( $c_{o}=1 / \sqrt{M C}$ ). Thus the Nyquist condition represents a computational bound, not a physical limitation. -

Problem \# 3: Now consider the cascade of $N$ such $\mathcal{T}(s)$ matrices and perform an eigenanalysis.

- 3.1: Find the eigenvalues and eigenvectors of $\mathcal{T}(s)$ as functions of $s / s_{c}$.

Sol: Matrix $\mathcal{T}(s)$ has eigenvalues

$$
\lambda_{ \pm}=1 \mp 2 s / s_{c} \approx e^{ \pm 2 s / s_{c}}=e^{\mp s T_{c}}
$$

From this we can interpret the eigenvalues as the cell delay $T_{c}=2 / s_{c}$.
The corresponding unnormalized eigenvectors are

$$
\boldsymbol{E}_{ \pm}=\left[\begin{array}{c}
\mp \sqrt{M / C} \\
1
\end{array}\right]
$$

where the characteristic impedance defined is $r_{o}=\sqrt{M / C}$.

## Problem \# 4: Find the velocity transferfunction $H_{12}(s)=V_{2} /\left.V_{1}\right|_{F_{2}=0}$.

-4.1: Assuming that $N=2$ and $F_{2}=0$ (two half-mass problem), find the transfer function $H(s) \equiv V_{2} / V_{1}$. From the results of the $\mathcal{T}$ matrix, find

$$
H_{21}(s)=\left.\frac{V_{2}}{V_{1}}\right|_{F_{2}=0}
$$

 $F_{2}=0$

$$
\frac{V_{2}}{V_{1}}=\frac{-1}{s^{2} M C / 2+1}=\left(\frac{c_{+}}{s-s_{+}}+\frac{c_{-}}{s-s_{-}}\right)
$$

having eigenfrequencies $s_{ \pm}= \pm \jmath \sqrt{\frac{2}{2 M C}}= \pm s_{c}$ and residues $c_{ \pm}= \pm \jmath / \sqrt{2 M C}= \pm s_{c}$.

$$
\text { - 4.2: Find } h_{21}(t) \leftrightarrow H_{21}(s)
$$

Sol:

$$
h(t)=\oint_{\sigma_{0}-j \infty}^{\sigma_{0}+j \infty} \frac{e^{s t}}{s^{2} M C / 2+1} \frac{d s}{2 \pi j}=c_{+} e^{-s_{+} t} u(t)+c_{-} e^{-s_{-} t} u(t)
$$

The integral follows from the Cauchy Residue theorem (CRT).

- 4.3: What is the input impedance $Z_{2}=F_{2} / V_{2}$, assuming $F_{3}=-r_{0} V_{3}$ ?

Sol: Starting from Eq. DE-3.4a find $Z_{2}$

$$
Z_{2}(s)=\frac{F_{2}}{V_{2}}=T\left[\begin{array}{c}
F \\
-V
\end{array}\right]_{2}=\frac{-\left(1+s^{2} C M / 2\right) r_{0} y_{2}-s M\left(1+s^{2} C M / 4\right) y_{2}}{-s C r_{0} y_{2}-\left(1+s^{2} C M / 2\right) y_{2}}
$$

■

- 4.4: Simplify the expression for $Z_{2}$ as follows:

1. Assuming the characteristic impedance $r_{0}=\sqrt{M / C}$,
2. terminate the system in $r_{0}: F_{2}=-r_{0} V_{2}$ (i.e., $-V_{2}$ cancels).
3. Assume higher-order frequency terms are less than $1\left(\left|s / s_{c}\right|<1\right)$.
4. Let the number of cells $N \rightarrow \infty$. Thus $\left|s / s_{c}\right|^{N}=0$.

When a transmission line is terminated in its characteristic impedance $r_{0}$, the input impedance $Z_{1}(s)=r_{0}$. Thus, when we simplify the expression for $\mathcal{T}(s)$, it should be equal to $r_{0}$. Show that this is true for this setup.

Sol: Applying the Nyquist approximation (i.e., ignore second order frequency terms $\left(s / s_{c}\right)^{2} \approx 0$ )

$$
\begin{aligned}
Z_{1}(s) & =\frac{r_{o}\left(1+\underline{s}^{2} C A / 2\right)^{0}+s M\left(1+\underline{s}^{2} C A / 4\right)^{0}}{r_{o} s C+\left(1+s^{2} C A / 2\right)^{0}} \\
& \approx \frac{r_{o}+s M}{1+r_{o} s C}=\frac{M C}{M C} \cdot \frac{r_{o}+s M}{1+r_{o} s C}=\frac{M}{C} \cdot \frac{r_{o} C+s M C}{M+r_{o} s M C}=r_{o}^{2} \frac{r_{o} C+s / s_{c}}{M+r_{o} s / s_{c}} \\
& \approx r_{o}^{2} \frac{r_{o} C+s \nmid \widehat{s}_{c}^{0}}{M+r_{o s \nmid s_{c}}^{0}}=r_{o}^{3} \frac{C}{M} \\
& =r_{o} .
\end{aligned}
$$

We conclude that below the Nyquist cutoff frequency, as $N \rightarrow \infty$ the system equals a transmission line terminated by its characteristic impedance thus $Z_{1}(s)=r_{o}$.

- 4.5: State the ABCD matrix relationship between the first and Nth nodes in terms of the cell matrix. Write out the transfer function for one cell, $H_{21}$.
Sol:

$$
\mathcal{T}=\left[\begin{array}{ll}
\mathcal{A} & \mathcal{B} \\
\mathcal{C} & \mathcal{D}
\end{array}\right]
$$

Now use the formulae for the eigenvalues and vectors to obtain $\mathcal{T}$ for $N=1$ :

$$
\mathcal{T}=E \Lambda E^{-1}=E\left[\begin{array}{cc}
\lambda_{+} & 0 \\
0 & \lambda_{-}
\end{array}\right] E^{-1}
$$

- 4.6: What is the velocity transfer function $H_{N 1}=\frac{V_{N}}{V_{1}}$ ?

Sol:

$$
\left[\begin{array}{l}
F_{1} \\
V_{1}
\end{array}\right]=\mathcal{T}^{N}\left[\begin{array}{c}
F_{N}(\omega) \\
-V_{N}(\omega)
\end{array}\right]
$$

along with the eigenvalue expansion

$$
\mathcal{T}^{N}=E \Lambda^{N} E^{-1}=E\left[\begin{array}{cc}
\lambda_{+}^{N} & 0 \\
0 & \lambda_{-}^{N}
\end{array}\right] E^{-1} .
$$

where $\lambda_{ \pm}^{N}=e^{\mp s N T_{o}}$. Recall that $N T_{o}$ is the one way delay.
We conclude that as we add more cells, the delay linearly increases with $N$, since each eigenvalue represents the delay of one cell, and delay adds.

